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Fundamental constants, conversions, etc.:

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Definite integrals:
\[ \int_{-\infty}^{\infty} e^{-x^2} \, dx = \sqrt{\pi}. \]  
(I–1)
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\[ \int \frac{1}{\sqrt{x^2 + a^2}} \, dx = \ln \left( x + \sqrt{x^2 + a^2} \right). \]  
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\[ \int \frac{1}{(x^2 + a^2)^{3/2}} \, dx = \frac{x}{a^2 \sqrt{x^2 + a^2}}. \]  
(I–5)

I–1. A magnetic field of magnitude \( B \) points in the direction \( \hat{z} \). The quantum-mechanical Hamiltonian for the spin states of an electron in this field is \( H = \omega S_z \), where \( \omega = \frac{|e|B}{meC} \), and \( S_z \) is the projection of the electron’s spin on the \( z \)-axis.

(a) Write down the time evolution operator which relates the states of the system at time \( t > 0 \) to the states at time \( t = 0 \) (in the Schrödinger picture).

(b) What are the stationary states of this system?

(c) If the initial state \( |\alpha, t_0 = 0\rangle \) is an eigenstate of the \( S_x \) operator (the spin projection along the \( \hat{x} \) axis) with eigenvalue \( \frac{1}{2} \hbar \) what is the state \( |\alpha, t > 0\rangle \)?

(d) Evaluate \( \langle \alpha, t | S_x | \alpha, t \rangle \), \( \langle \alpha, t | S_y | \alpha, t \rangle \) and \( \langle \alpha, t | S_z | \alpha, t \rangle \) and argue that the spin precesses around the \( \hat{z} \) axis.

I–2. (a) Consider a non-relativistic particle of mass \( m \) in a spherically symmetric classical potential energy \( V(r) \), where \( r \) is the distance of the particle from a fixed point. Derive an equation for the radial motion of the form
\[ m \frac{d^2 r}{dt^2} = -\frac{dV_{\text{eff}}(r; L)}{dr}, \]  
(I–6)
where the \( V_{\text{eff}} \) is an effective potential for radial motion, and depends on the angular momentum \( L \).
(b) In general relativity, it is found that an equation of this form still applies to the motion of a small object around another heavy spherical object of mass $M$. But the effective potential energy is now

$$V_{\text{eff}}(r; L) = \frac{m}{2} \left( 1 - \frac{2GM}{c^2 r} \right) \left( \frac{L^2}{m^2 r^2} + c^2 \right). \quad (I-7)$$

Show that in the limit $c \to \infty$, this is equivalent to the situation for Newtonian gravity.

The remaining parts of this problem refer to motion governed by Eq. (I-7).

(c) For which values of $L$ do circular orbits exist when the effective potential energy is (I-7)? Find the radii of stable circular orbits and their angular velocities $\omega_\phi$.

(d) If an orbiting object is slightly disturbed from a stable circular orbit at radius $R$, it oscillates around the stable radius. Compute the corresponding oscillation frequency $\omega_r$. Expand the precession rate $\omega_p := \omega_\phi - \omega_r$ for $R \gg GM/c^2$ to leading non-vanishing order.

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I–3. A hollow, empty, irregularly shaped conductor is charged to a potential $V$ relative to infinity. Someone has calculated the electrostatic potential $\phi$ for the system, and far from the conductor the result is

$$\phi(r) = \frac{AV}{r} + \frac{BV z}{r^3} + \frac{CV}{r^3} \left( \frac{3 z^2}{r^2} - 1 \right) \quad (I-8)$$

where $r^2$ is $|r|^2 = x^2 + y^2 + z^2$, and $A$, $B$, and $C$ are constants.

(a) What is leading-order contribution to the electric field at large distances from the conductor?

(b) What is the charge on the conductor?

(c) What is the capacitance of the conductor?

---

I–4. A monatomic classical gas of molecular weight $M$ has a temperature $T$ at low pressure. There are $N$ molecules in a volume $V$. The velocity distribution function for the velocity of the molecules is

$$f(v) = K \exp[-\beta(v_x^2 + v_y^2 + v_z^2)], \quad (I-9)$$

where $K$ is a constant. (Thus the number of molecules in a volume element $d^3v$ is $f(v) d^3v$.)

(a) What is the value of $\beta$?
(b) Derive an equation for the distribution of the speed of the molecules, i.e. a function $g(v)$ such that $g(v)dv$ is the probability of finding a particle with speed $v$ to $v + dv$.

(c) Now let a small hole of area $A$ be made in a container of such a gas, to give an effusive molecular beam. Obtain an analytic formula for the rate at which molecules leave the container. Explain why the distribution of molecular speeds emerging from the hole has a different shape to that you found in the previous part.
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II–1. A copper wire of diameter \( d = 1 \) mm is shaped in the form of a vertical square loop with sides of length \( a = 10 \) cm. It falls with velocity \( v \) from a region where there is a constant horizontal magnetic field of magnitude 1.2 T perpendicular to the plane of the loop, into a region with zero magnetic field, as shown in the figure below. The resistivity of the wire is \( \rho_e = 1.7 \times 10^{-8} \) \( \Omega \cdot \text{m} \), and its density is \( \rho_m = 8960 \) kg/m\(^3\).

![Figure II–1: For problem II–1.](image)

(a) In what direction does a current flow around this loop, and what is its magnitude in terms of the velocity of fall, \( v \)?

(b) Show that this gives an upward force on the loop.

(c) Find an expression for the terminal velocity. (I.e., the velocity at which the acceleration would be zero, while the boundary between the two regions of field intersects the circuit.) Estimate the value of the terminal velocity, taking \( g = 9.8 \) m/s\(^2\).
II–2. In this problem you will investigate the conditions limiting the braking ability of a bicycle.

The bicycle moves on a horizontal road in a straight line. See Fig. II–2 for the geometry. The distance between the points of contact of the wheels is 1.05 m. The center of gravity of the bicycle-plus-rider system is 1.15 m above the ground and 0.40 m in front of the point of contact of the rear wheel. In the following calculations, ignore the rotational inertia of the wheels.

There is a coefficient of friction $\mu = 0.8$ between each of the tires and the road.

(a) When only the brake on the rear wheel is applied, what is the maximum deceleration? What causes the limit?

(b) Repeat when only the front brake is used.

II–3. In the variational method for estimating the solution to a quantum-mechanical bound-state problem, show that a variational estimate for the ground-state energy is equal to or larger than the true ground-state energy.

Apply the variational method to a particle of mass $m$ in a spherically symmetric Yukawa potential

$$ V(r) = \frac{g^2 e^{-r\mu}}{r}, \quad (II–6) $$

where $r$ is the distance between the particle and the fixed center of force. Use a trial wave function of the form

$$ \psi(r) \propto e^{-r/a}. \quad (II–7) $$

Here $a$ is a parameter to be determined. It will be sufficient to carry through your calculation to where $a$ and the ground state energy can be determined by an algebraic calculation.
II–4. Consider an ideal gas (of 1 mole) that obeys the usual equation \( PV = RT \). The internal energy depends on \( T \), and is independent of \( P \) and \( V \).

(a) Starting from the first law of thermodynamics, show that

\[
C_P - C_V = R.
\]  

(II–8)

(b) Show that for a quasistatic adiabatic process

\[
PV^\gamma = \text{constant},
\]  

(II–9)

where \( \gamma = C_P/C_V \). (Assume the specific heats are constant.)

(c) A sketch of the \( P-V \) diagram of an isothermal expansion from an initial state \( i \) to a final state \( f \) is shown below. Sketch the curve of an adiabatic process starting from the same initial state. Does the temperature rise or fall in the process?